

# Thermoelectric cooling properties of a quantum Hall Corbino device

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We analyze the thermoelectric cooling properties of a Corbino device in the quantum Hall regime on the basis of experimental data of electrical conductance. We focus on the cooling power and the coefficient of performance within and beyond linear response. Thermovoltage measurements in this device reported in *Phys. Rev. Applied*, **14** 034019 (2020) indicated that the transport takes place in the diffusive regime, without signatures of effects due to the electron-phonon interaction in a wide range of temperatures and filling factors. In this regime, the heat and charge currents by electrons can be described by a single transmission function. We infer this function from experimental data of conductance measurements and we calculate the cooling power and the coefficient of performance for a wide range of filling factors and temperatures, as functions of the thermal and electrical biases. We predict an interesting cooling performance in several parameter regimes.

## I. INTRODUCTION

Thermoelectric phenomena in quantum mesoscopic devices is a topic of great interest for some years now[1–4]. The relevant regime takes place at very low temperatures, typically in the subkelvin range. Besides the interest in the fundamental physical properties, research in this field is motivated by the emergent development of quantum technologies. For instance, all the proposals for implementing quantum computation in solid-state devices rely on systems that operate at cryogenic temperatures, being one of such platforms the fractional quantum Hall system[5, 6]. One of the challenges in this context is the removal of heat produced in the operation of these devices.

A significant amount of work has been published on thermoelectricity in quantum coherent systems. Paradigmatic examples are quantum dots[7–16], nanowires[17, 18] quantum point contacts[19–22] and topological edge states of the quantum Hall and the quantum spin Hall regimes[23–35]. In these cases transport takes place through a few ballistic channels. In two-terminal configurations, for very low temperatures, such that electron-phonon interaction does not play a relevant role, charge and thermal currents, as well as the thermoelectric response are fully characterized by a transmission function that describes the transport properties of the electrons through the device [1, 20, 21, 36–39].

Thermoelectricity in the quantum Hall regime in the Corbino geometry has been investigated theoretically[40, 41] and experimentally[42, 43], as well as in bar configurations[44, 45]. In a recent work, experimental evidence of a significant thermoelectric response in a quantum Hall Corbino device was presented [46]. The theoretical modeling of this experimental data is consistent with a diffusive mechanism of the electron transport, free from

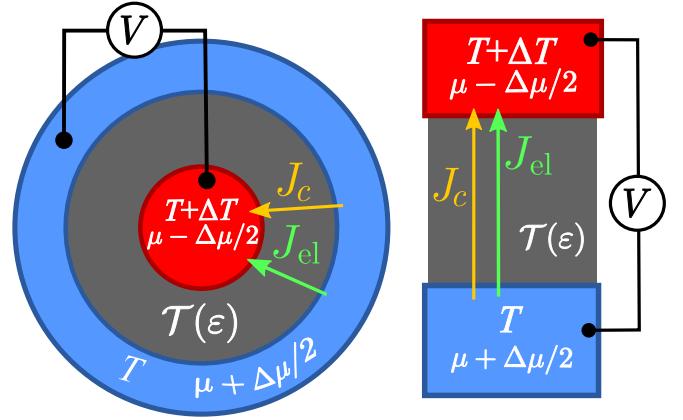


FIG. 1. Left panel: sketch of the Corbino ring. The spectrum of Landau levels define a conductor characterized by a transmission function  $\mathcal{T}(\varepsilon)$  between two biased reservoirs with a temperature difference  $\Delta T$ . Applying an external voltage  $V$  a chemical potential difference induced, leading to a cooling thermal current  $J_c$  and an electrical current  $J_{el}$ . Right panel: equivalent two terminal configuration, characterized by the same transmission probability and identical temperature and voltage biases.

the effects of electron-phonon interaction, accounting for the thermoelectric response within partially filled Landau levels. This description is based on a single transmission function, which determines all the transport coefficients, as in the case of the ballistic regime.

A necessary ingredient for the thermoelectric heat-work conversion is the existence of energy filters, implying spectral properties of the device in which the transport of electrons and holes is asymmetric. This implies a transmission probability having a profile with rapid changes as a function of the energy. The mathemati-

cal properties that a transmission function must satisfy for an efficient thermoelectric performance were analyzed in several works [20, 21, 33, 35, 38, 47].

The aim of the present contribution is to present a detailed theoretical study of the cooling properties of the Corbino device investigated in Ref. 46 on the basis of the transmission function inferred from the experimental data for the conductance recorded in that work. A sketch of the considered setup is shown in Fig. 1. The two-dimensional electron system threaded by an external high magnetic field is constructed with a circular shape, hosting a central hot region, as indicated in the left panel. The system can be electrically biased by means of a voltage difference. The electrons in the quantum Hall regime are accommodated in Landau levels separated by a gap, which play the role of energy filters defining the thermoelectric performance. The possibility of realizing thermoelectric cooling with the Corbino geometry was previously suggested in Ref. [48], where the term ‘‘Landau cooling’’ has been coined. The device effectively behaves as a two-terminal setup characterized by a transmission function, as indicated in the right panel of the Figure. The appealing of using the Landau levels for the thermoelectric heat–work conversion mechanism is the simplicity of its implementation in comparison to using edge states, which requires more complex structures with quantum point contacts and quantum dots.

The work is organized as follows. In section II we present the procedure to infer the transmission function from the data of the conductance and we define the theoretical treatment. In section III we analyze the regime for which cooling is possible. The performance is analyzed in Section IV and Section V is devoted to summary and conclusions.

## II. THEORETICAL APPROACH

### A. Inference of the transmission function from experimental data

In Ref. [46] measurements for the conductance and the thermovoltage of a Corbino device with a central heater were reported. The experimental data was successfully explained on the basis of calculations of the Onsager transport coefficients in terms of a transmission function, which was inferred from the conductance. Here, we will also infer the transmission function following a similar procedure in order to evaluate the cooling power. We summarize the steps of Ref. [46] to infer the transmission function of the Corbino device from the knowledge of the experimental data for the conductance,  $G_{\text{exp}}(B, T_0)$ , as a function of the magnetic field  $B$  at a low temperature  $T_0$ .

In the setup sketched in the left panel of Fig. 1, the conductance is obtained by measuring the electric current as a function of the magnetic field with the inner and outer reservoirs at the same temperature  $T_0$ . The sys-

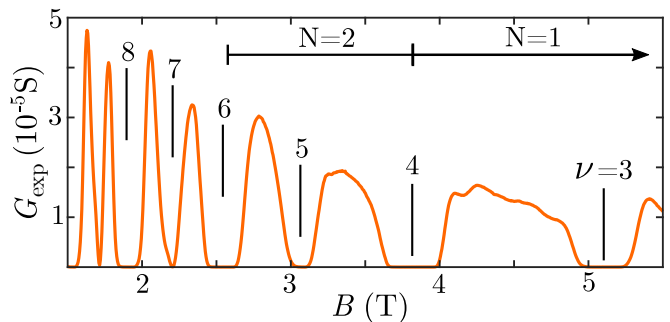


FIG. 2. Conductance measurement of a Corbino device as a function of the magnetic field for  $T_0 = 269$  mK. Filling fractions  $\nu$  and Landau levels  $N$  are indicated.

In the text, we will refer to the ‘‘peak  $\nu$ ’’ for ‘‘the peak before the filling factor  $\nu$ ’’. The measurement shown corresponds to a Corbino sample of 3.5 mm in diameter, having several concentric ohmic contacts and a central heater. The particular Corbino ring measured in this figure being of 600  $\mu\text{m}$  and 1600  $\mu\text{m}$  inner and outer ring respectively. The sample was grown by molecular beam-epitaxy on a GaAs wafer having a single 2DES. From van-der-Pauw geometry measurements on test pieces a mobility of  $21 \times 10^6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and an electron density of  $2.0 \times 10^{11} \text{ cm}^{-2}$  was determined at 1.3 K in the dark [46].

tem is biased with a small voltage difference  $V$  applied radially. The experimental conductance  $G_{\text{exp}}(B, T_0)$  is directly given by the ratio between this current and  $V$ . As indicated in the right panel of the Figure, the setup can be effectively regarded as a two-terminal configuration, where the transport properties are described by the transmission probability  $\mathcal{T}(\varepsilon)$ . Formally, the corresponding conductance at a temperature  $T$  is related to this function through

$$G(\mu, T) = -\frac{e^2}{h} \int_{-\infty}^{+\infty} d\varepsilon \mathcal{T}(\varepsilon) \frac{\partial f(\varepsilon)}{\partial \varepsilon}, \quad (1)$$

where  $f(\varepsilon) = 1/(e^{(\varepsilon-\mu)/k_B T} + 1)$  is the Fermi-Dirac distribution function, which depends on the reference chemical potential  $\mu$  and the temperature  $T$ . Deep in the quantum Hall regime, where the contributions of the different integer filling fractions are clearly separated one another, the Fermi energy is related to the magnetic field  $B_\nu = n_e h/e\nu$  corresponding to the filling fraction  $\nu$  as  $\varepsilon_F = \hbar e B_\nu / 2m^*$ , being  $n_e$  the electron density of the sample. In a real sample, Landau levels are broadened and the previous relations extend to a range of magnetic fields around  $B_\nu$ . On the other hand, for low enough temperatures, the derivative of the Fermi function can be approximated as  $-\partial f(\varepsilon)/\partial \varepsilon \rightarrow \delta(\varepsilon - \mu)$ . Hence, Eq. (1) can be expressed as  $G(\mu, T) = (e^2/h) \mathcal{T}(\mu)$ , while  $\mu = \varepsilon_F$ . Therefore, assuming that the temperature  $T_0$  is low enough, we can calculate the transmission function from the experimental conductance  $G_{\text{exp}}(B, T_0)$  as follows,

$$\mathcal{T}(\varepsilon) = \frac{h}{e^2} G_{\text{exp}}(B, T_0), \quad (2)$$

with  $\varepsilon = \hbar e B / 2m^*$ . Experimental conductance results at a temperature  $T_0 = 269$  mK are shown in Fig. 2. We can see in the figure a profile of several features separated by an energy gap. Each of these features is associated to a filled Landau level (identified with a label  $N$ ), which is widened as a consequence of disorder in the sample and split because of the Zeeman interaction. The corresponding integer filling factors are indicated with  $\nu$  in the figure. Notice that while this inference of the transmission function is an exact procedure for a conductance recorded at  $T_0 = 0$ , there is a distortion as a consequence of the finite temperature at which the experiment is performed, which we will neglect in what follows. The alignment of the Landau levels to the reservoirs chemical potential is modified by means of the magnetic field. Hence, it can be regarded as a knob to tune the transport properties of the device.

This profile contains features similar to smoothed step functions separated by wells and also Lorentzian-type peaks. Since the cooling properties of transmission functions with these characteristics were recently carefully analyzed in Ref. [35], we will take into account several insights from that work as a reference for our study.

## B. Cooling current and electrical power

We now consider the situation where, in addition to a voltage bias  $V$ , a temperature bias associated to a thermal load, is radially applied. In order to focus on a configuration close to the experimental one studied in Ref. [46], which contained a heater in the center of the sample, we present in Fig. 1 the hottest reservoir at the center. However, we stress that the forthcoming analysis does not depend on whether the hot reservoir is placed in the central region of the Corbino or in the external rim. Furthermore, we highlight the fact that the present configuration can be actually regarded as a two-terminal setup under the effect of a thermal bias  $\Delta T$ , and an electrical bias  $\Delta\mu = eV$ , as indicated in the right-hand side of Fig. 1. In this representation it is clear that the position of the cold and hot reservoirs can be exchanged.

Following Ref. [35], we focus on the heat current transported by the electrons, exiting the reservoir at the lowest temperature  $T$ . In terms of the transmission function, it reads

$$J_q = -\frac{1}{h} \int_{-\infty}^{+\infty} d\varepsilon \mathcal{T}(\varepsilon) \left( \varepsilon - \mu - \frac{\Delta\mu}{2} \right) \Delta f(\varepsilon), \quad (3)$$

being  $\Delta f(\varepsilon) = f_c(\varepsilon) - f_h(\varepsilon)$ , the difference between the Fermi functions corresponding to the cold/hot (c/h) reservoir at temperatures  $T_c = T$ ,  $T_h = T + \Delta T$  and chemical potentials  $\mu_c = \mu + \Delta\mu/2$  and  $\mu_h = \mu - \Delta\mu/2$ , respectively. The heat flux given by Eq. (3) is directly proportional to the entropy flux entering or exiting the coldest reservoir due to the particle flux through the device. In the absence of inelastic scattering processes, the

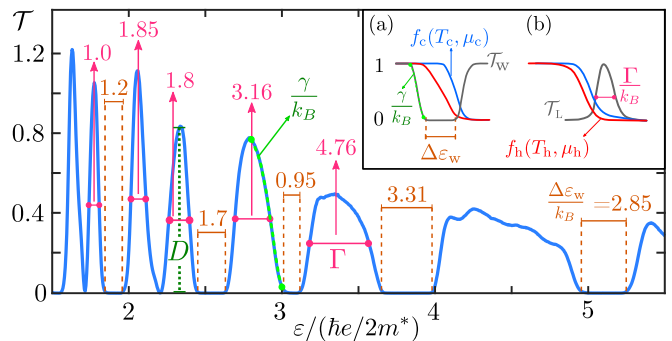


FIG. 3. Transmission function inferred from the measured conductance shown in Fig. 2. Equivalent temperature estimates in kelvin for the gaps (orange), half height width of the Landau levels  $\Gamma/k_B$  (pink), smoothing parameter  $\gamma/k_B$  and associated height parameter  $D$  are indicated. These will impact the cooling range as well as the the optimal operational conditions (see text).

*Inset:* (a) Well-like transmission function  $\mathcal{T}_w$  (grey) including the Fermi functions  $f_h$  for the hot (red) and the cold  $f_c$  (blue) reservoirs. The width of the well is  $\varepsilon_w$ . (b) Lorentzian-like transmission function  $\mathcal{T}_L$  of width  $\Gamma$  and height  $D = 1$ .

latter is fully characterized by its transmission probability  $\mathcal{T}(\varepsilon)$ .

The heat flux described by Eq. (3) corresponds to a cooling process when it exits the coldest reservoir. We define the *cooling power* as

$$J_c \equiv J_q, \text{ if } J_q \geq 0, \quad J_c \equiv 0, \text{ otherwise.} \quad (4)$$

Similarly to other works, we will consider as a reference the Pendry quantum bound[20, 21, 49]

$$J_{qb} = \frac{\pi^2 k_B^2}{6h} T^2, \quad (5)$$

which defines the maximum heat current that can be transported through a quantum channel of unitary transmission.

The thermal and electrical biases induce an electron flux  $J_{el}$ , and have an associated electrical power  $P$ , which are given by

$$J_{el} = -\frac{1}{h} \int_{-\infty}^{+\infty} d\varepsilon \mathcal{T}(\varepsilon) \Delta f(\varepsilon), \quad P = \Delta\mu J_{el}. \quad (6)$$

## C. Cooling performance

It is natural to quantify the efficiency of the cooling operation in terms of the coefficient of performance, defined as

$$\text{COP} = \frac{J_c}{P}, \quad (7)$$

with the definitions of the cooling power and the invested electrical power given by Eqs. (4) and (6), respectively. This quantity is bounded by the Carnot limit  $\eta_C = T/\Delta T$ .

A particular important regime corresponds to small amplitudes of the thermal and electrical biases, where  $\Delta T/T$  and  $\Delta\mu/k_B T$  are small enough to justify an expansion up to linear order in these quantities in the expressions of the heat and charge currents defined, respectively, in Eqs. (3) and (6). The result is

$$\begin{pmatrix} J_{\text{el}}/e \\ J_{\text{q}} \end{pmatrix} = \begin{pmatrix} \mathcal{L}_{11} & \mathcal{L}_{12} \\ \mathcal{L}_{21} & \mathcal{L}_{22} \end{pmatrix} \begin{pmatrix} \Delta\mu/k_B T \\ \Delta T/k_B T^2 \end{pmatrix}. \quad (8)$$

where  $\hat{\mathcal{L}}$  is the Onsager matrix, with elements calculated from the transmission function as follows

$$\mathcal{L}_{ij} = -T \int \frac{d\varepsilon}{h} \frac{\partial f(\varepsilon)}{\partial \varepsilon} (\varepsilon - \mu)^{i+j-2} \mathcal{T}(\varepsilon). \quad (9)$$

Notice that  $\mathcal{L}_{12} = \mathcal{L}_{21}$ , since the present device effectively behaves like a usual two-terminal one, as sketched in Fig. 1.

In this regime, the quantity characterizing the maximum COP for a fixed value of  $\Delta T$  is parametrized by the figure of merit  $ZT = \mathcal{L}_{21}^2 / \text{Det} \hat{\mathcal{L}}$  through[1]

$$\text{COP} = \eta_C \frac{\sqrt{ZT + 1} - 1}{\sqrt{ZT + 1} + 1}. \quad (10)$$

Hence, the ideal Carnot limit is achieved for  $ZT \rightarrow \infty$ .

The usual situation in thermoelectricity is that devices for which the COP is high, the cooling power is low and viceversa [1, 20, 21, 33, 35, 38, 47]. Therefore, we will analyze separately the following aspects: (i) under which conditions is cooling possible in the device, (ii) the conditions for which a large cooling power is expected and (iii) the conditions for a large COP.

#### D. Conditions for cooling

Before going to the concrete analysis of the cooling properties of the sample for which the conductance is reported in Fig. 2, it is useful to summarize the outcome of Ref. [35] regarding the properties of  $\mathcal{T}(\varepsilon)$  leading to cooling in transmission functions containing peaks, steps and well-type features.

The optimal profile leading to the highest cooling power is a single rigid step function of the form  $\mathcal{T}_{\text{step}}(\varepsilon) = D\theta(\varepsilon - \varepsilon_w)$ . The optimal operational condition corresponds to the chemical potential of the coldest reservoir coinciding with the onset of the step,  $\mu_c = \varepsilon_w$ , and large electrical bias, with  $\mu_h$  as far apart as possible from the step. This configuration of chemical potentials enables the thermoelectric compensation of the effect of the heat flow because of the temperature bias, leading to the maximum bound for the cooling power,  $J_c^{\text{max}} = DJ_{\text{qb}}/2$ .

If, instead of a sharp step, the transmission function has the shape of a smoothed one characterized by a smoothness parameter  $\gamma$ , the maximum achievable cooling power decreases in reference to  $J_c^{\text{max}}$  by a factor  $DJ_{\text{qb}}(\gamma/k_B T)^2/4$ , which can be achieved for  $T > \gamma/4k_B$

and similar configurations of the chemical potentials as in the case of a sharp step. Furthermore, for temperatures  $T < \gamma/4k_B$ , cooling is strongly depressed and it is possible only for small  $\Delta\mu$ .

On the other hand, the cooling operation becomes strongly limited at high temperatures by the existence of well-type features. A sketch of such transmission function, denoted by  $\mathcal{T}_w(\varepsilon)$  and characterized by a typical energy  $\Delta\varepsilon_w$ , is indicated in the inset (1) of Fig. 3, along with the two Fermi functions  $f_c(\varepsilon)$  and  $f_h(\varepsilon)$ . In such cases, starting from a situation where the chemical potential  $\mu_c$  is fixed at one of the step functions, the amplitude of the bias voltage –applied to compensate the thermal bias– is limited by the existence of the mirror step function. Hence, the cooling operation in this case is limited by  $T_h$ ,  $T_c \sim \Delta\varepsilon_w/k_B$ . The inset of Fig. 3 shows a configuration with the chemical potential of the cold reservoir selected at the step with positive slope, in which case  $\Delta\mu = \mu_c - \mu_h > 0$ . An identical result can be obtained when  $\mu_c$  is placed at the step with negative slope (the mirror-related step) and  $\Delta\mu = \mu_c - \mu_h < 0$ .

In the case of sharp peaks (Lorentzian-like) features, there are other limiting properties in the operation. An example of such a transmission function is sketched in the inset (b) of Fig. 3, denoted by  $\mathcal{T}_L(\varepsilon)$  and characterized by a width  $\Gamma$ . In the present case, cooling can also be achieved for  $\mu_c$  placed at the positive slope of the peak and  $\Delta\mu = \mu_c - \mu_h > 0$  applied to compensate the thermal bias or, equivalently, when  $\mu_c$  is placed at the negative slope of the peak and  $\Delta\mu = \mu_c - \mu_h < 0$  applied. The cooling operation is limited in this case to the regime where  $T_c, T_h < \Gamma/4k_B$  and is possible only for small values of  $|\Delta\mu|$ .

### III. COOLING POWER

In Fig. 3, the parameters characterizing the different features of the transmission function  $\mathcal{T}(\varepsilon)$  are indicated. In particular, widths of the main peaks and wells, respectively  $\Gamma$  and  $\Delta\varepsilon_w$  are indicated in units of temperature in order to facilitate the comparison of these properties with the thermal width entering the function  $\Delta f$  in Eq. (4), which we consider to be  $\sim 4k_B T$ . The height  $D$  of each feature can be simply read in the vertical axis.

At high temperatures, the main limitation for the cooling range is the width of gaps between the different Landau levels. If the gap between two consecutive Landau levels is of the same order of magnitude of the transport window, cooling is no longer possible.

The cooling power normalized by the quantum bound,  $J_c/J_{\text{qb}}$ , as a function of  $\mu$  and  $\Delta\mu$  in a range of temperatures  $200 \text{ mK} < T < 800 \text{ mK}$  is shown in Fig. 4 for a fixed ratio  $\Delta T/T_h = 0.1$ . In this Fig. we can identify different regions of values of chemical potentials and voltages for which cooling is possible in the device within this temperature regime. We see that they are localized close to values of energies for which  $\mathcal{T}(\varepsilon)$  has rapid changes and

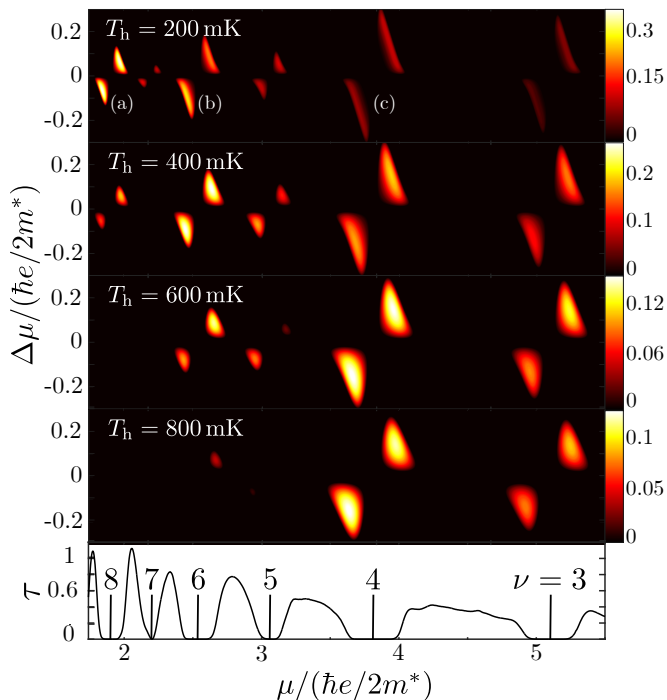


FIG. 4. Contour plots for the cooling power normalized by the quantum bound,  $J_c/J_{qb}$  as a function of  $\mu/(\hbar e/2m^*)$  and  $\Delta\mu/(\hbar e/2m^*)$  at a fixed relation  $\Delta T/T_h = 0.1$ .  $\mathcal{T}(\varepsilon)$  within the same range of energies is shown as a reference in the bottom panel. Note that the color scale changes on each panel.

the profile of  $\mathcal{T}(\varepsilon)$  is similar to well-type  $\mathcal{T}_w$  and peak-like  $\mathcal{T}_L$  transmission functions discussed before. This takes place close to the onset and closing of the different Landau levels. As mentioned in Section IID, the sign of  $\Delta\mu$  leading to cooling depends on the sign of the corresponding slope of  $\mathcal{T}(\mu_c)$ . For the lowest temperature, shown in the upper panel, we see that the highest intensity corresponds to the peaks  $\nu = 6, 7$  and  $8$ . This is because in this temperature range,  $T_c$  and  $T_h$  are lower than the width of these peaks, as well as lower than the gaps between the peaks  $\nu = 7$  and  $8$  and between  $\nu = 5$  and  $6$ . Under these conditions, the features favoring the strength of the power are the slope ( $1/\gamma$ ) and the height ( $D$ ) of  $\mathcal{T}(\varepsilon)$ , and these two peaks are the optimal ones regarding these properties. The estimates are  $\gamma/4k_B \sim 360$  mK,  $329$  mK and  $190$  mK for  $\nu = 6, 7, 8$ , respectively, while  $D \sim 1.2$ . On the other hand, as the temperature grows (see top to bottom panels in the Fig.), the cooling lobes centered on these filling factors become fainter and disappear when the temperature becomes comparable to the width of the gap between neighboring peaks. In the high-temperature range, the optimal cooling region is shifted towards lower filling factors ( $\nu = 3, 4$ ). The cooling power is decreased in comparison to the previous cases since the slopes and heights are smaller in these cases. However, the range of  $\mu, \Delta\mu$  for which cooling is possible is larger for these lower filling factors, since the widths  $\Gamma$  of the Landau levels and the gaps  $\Delta\varepsilon_w$  become larger.

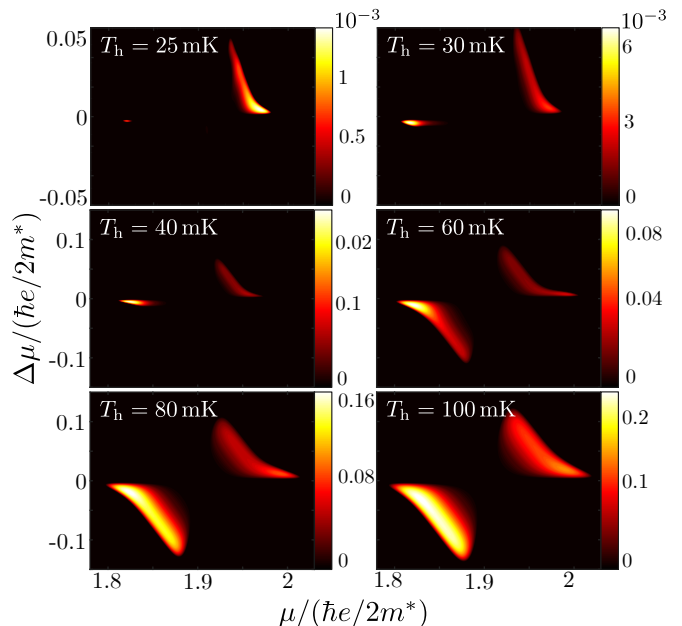


FIG. 5. Similar to Fig. 4 with focus on the low temperature regime. Contour plots corresponding to the vicinity of  $\nu = 8$ .

A similar analysis with focus on the low-temperature regime is presented in Fig. 5. These plots show the cooling power within a range of temperature  $25 \text{ mK} < T < 100 \text{ mK}$ . The optimal range of magnetic fields (or equivalent chemical potentials) in this case corresponds to the one indicated with (a) in Fig. 4 and is centered between  $\nu = 7$  and  $8$ . These peaks of  $\mathcal{T}(\varepsilon)$  have approximately the same width  $\Gamma/4k_B \sim 460$  mK and a corresponding slope  $\gamma/4k_B \sim 360$  mK. The range of temperatures covered in Fig. 5 have an associated thermal broadening,  $\sim 4k_B T_c$ , which is smaller than  $\gamma$ . This explains the low values of the cooling power, which shows a decreasing behavior as the temperature decreases. Notice that the highest values of  $J_c$  are achieved in this region of chemical potentials close to  $T_h \sim 200$  mK (shown in the top panel of Fig. 4), which corresponds precisely to the regime where  $T_c$  becomes larger than  $\gamma/4k_B$ .

We close the analysis of the cooling power with the results shown in Fig. 6. These correspond to the cooling power as a function of the temperature of the hot bath  $T_h$  and the relative temperature difference  $\Delta T/T_h$  for fixed values of  $\mu$  and  $\Delta\mu$ . These were chosen to coincide with the ones identified with (a), (b) and (c) in Fig. 4 and are the ones leading to the highest cooling power close to  $\nu = 8, 6, 4$ , respectively. From the information of this Fig. we see that cooling at low temperatures is possible only near  $\nu = 8$ , corresponding to panel (a). This is precisely the case analyzed in Fig. 5. Similar values of the cooling power can be achieved for the parameters analyzed in panel (b). In this case, cooling is only possible above  $T \sim 100$  mK. Hence, the low temperature regime analyzed in Fig. 5 cannot be reached by cooling with the peak  $\nu = 6$ . The reason is that the transmission function

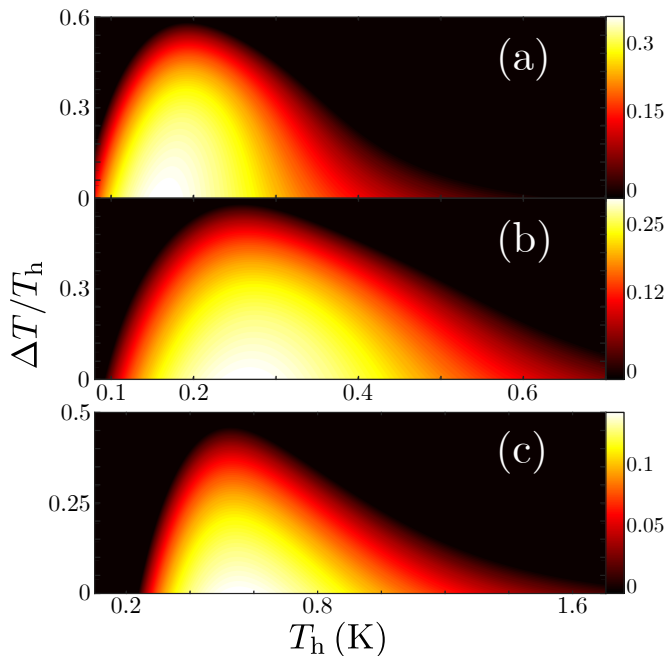


FIG. 6. Contour plots for the cooling power normalized by the quantum bound. The fixed values of  $\mu$  and  $\Delta\mu$  are marked in Fig. 4 and correspond (in units of  $\hbar e/2m^*$ ) to  $\mu_{(a)} = 1.85$ ,  $\Delta\mu_{(a)} = -0.064$ ;  $\mu_{(b)} = 2.459$ ,  $\Delta\mu_{(b)} = -0.095$  and  $\mu_{(c)} = 3.654$ ,  $\Delta\mu_{(c)} = -0.144$ .

peak in the latter case is smoother than in the case (a) (compare  $\gamma/4k_B \sim 625$  mK to  $\gamma/4k_B = 190$  mK of the previous case). On the other hand, cooling is possible up to higher temperatures in case (b), up to  $T_h \sim 750$  mK. Since  $\Gamma/4k_B$  as well as  $\Delta\varepsilon_w$  are larger in (b) than in (a). Similar observations can be done on the cooling power of the Landau level peak below  $\nu = 4$ , which is shown in panel (c). In that case, the range of temperature for cooling is the largest one, as a consequence of the large values of  $\Gamma$  and  $\Delta\varepsilon_w$  associated to this Landau level. The onset for cooling is at a higher temperature than in the previous cases,  $T_h \sim 200$  mK as a consequence of the larger smoothness  $\gamma/4k_B \sim 720$  mK. The largest achieved values for the cooling power are a bit lower than in the previous cases, as a consequence of the smaller height  $D$  of  $\mathcal{T}(\varepsilon)$ .

The absolute maximums of the cooling power as a function of temperature are determined by the smoothness parameter  $\gamma$  of the different Landau levels and the width  $\varepsilon_w$  of the gap. Therefore, they take place at  $\nu = 8$  for temperatures below 300 mK, at  $\nu = 6$  and at  $\nu = 4$  for temperatures above  $\sim 600$  mK. It is also interesting to note that the specific values of  $\mu$  and  $\Delta\mu$  do not change considerably and the temperature changes. For example if we focus on the lobe (a) in Fig. 4 the cooling area decreases as the temperature increases. However, the values of  $\mu$ ,  $\Delta\mu$  corresponding to the maximum position remain almost the same until the temperature is high enough that the maximum cooling power takes place at

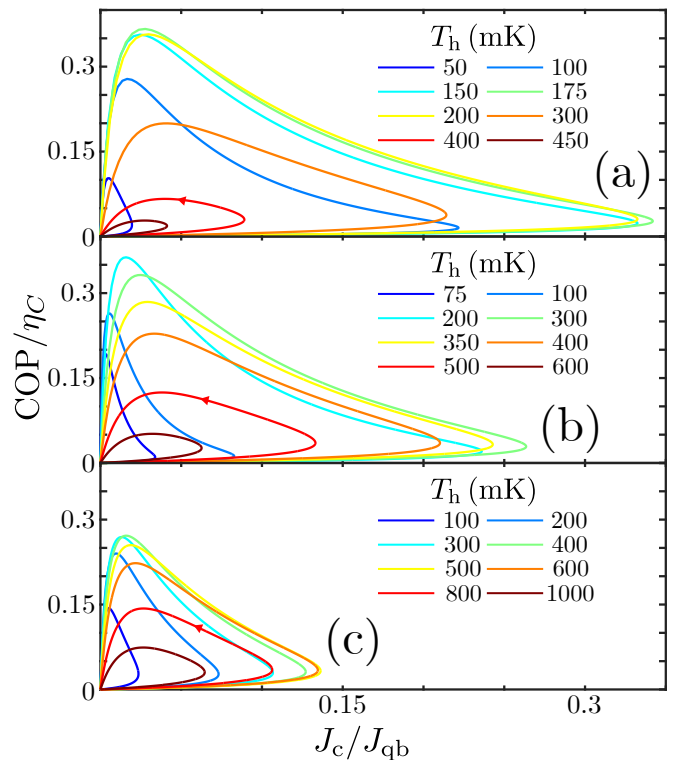


FIG. 7. Coefficient of performance normalized by the Carnot efficiency ( $\eta_C = T_c/\Delta T$ ) versus the cooling power normalized by the quantum bound ( $J_{qb}$ ) for varying  $\Delta\mu/(\hbar e/2m^*)$  from  $-0.25$  to  $0.25$ . The values of  $\mu/(\hbar e/2m^*)$  are 1.844, 2.459 and 3.654, corresponding to points (a), (b) and (c) indicated in Fig. 4, respectively.  $\Delta T/T_h = 0.1$  in all cases.

the filling factor, indicated with (b). This happens at a temperature,  $\sim 400$  mK. For even higher temperatures  $\sim 600$  mK, a similar effect takes place between the lobes associated to (b) and (c).

#### IV. COOLING EFFICIENCY

The usual situation in thermoelectric devices is that the highest efficiencies are achieved for parameters where the cooling power is very low and vice-versa and this is also the case in the present device. For this reason, it is very useful to represent both quantities at the same time in the form of the so called “lasso” plots, where COP vs  $J_c$  are represented for a set of parameters.

Fig. 7 shows “lasso” representations of COP vs  $J_c$  as  $\Delta\mu/(\hbar e/2m^*)$  changes within a specific interval while the other parameters correspond to the ones indicated with (a) (b) and (c) in Fig. 4 as (a), (b), (c). As mentioned before, the largest values of  $J_c$  correspond to smallest values of COP and vice-versa, but we can see wide ranges of parameters where both qualifiers achieve interesting values.

So far, we have not introduced any assumption regarding the amplitudes of the thermal and electrical bi-

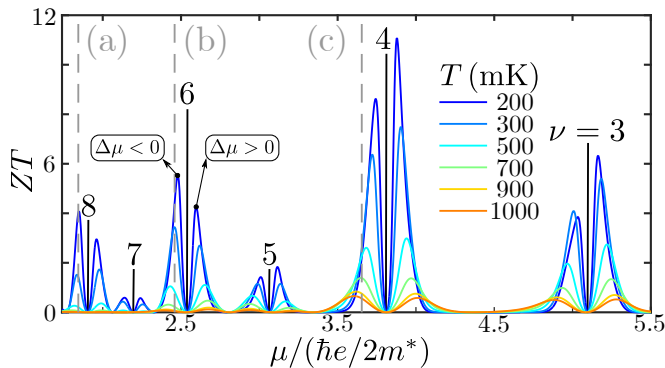


FIG. 8. Figure of merit,  $ZT$  intercepted by the  $\mu/(\hbar e/2m^*)$  values of points (a), (b) and (c) as indicated in Fig. 4. We note that the highest performance is in the vicinity of the fill factor  $\nu = 4$ .

ases. For the case where  $\Delta T/T_c$  and  $\Delta\mu/k_B T_c$  are small enough to justify an expansion up to linear order as in Eq. (8) the COP is given by Eq. (10).

The corresponding results are shown in Fig. 8. In order to get cooling, it is necessary to apply a voltage bias with a positive or negative sign, depending on the properties of  $\mathcal{T}(\varepsilon)$ , which define the behavior of  $\mathcal{L}_{ij}$ . Crucially, the sign of the off-diagonal coefficients,  $\mathcal{L}_{12} = \mathcal{L}_{21}$  depend on the sign of the slope of  $\mathcal{T}(\mu_c)$ . Thus,  $\Delta\mu > 0$  ( $\Delta\mu < 0$ ) for positive (negative) slope of  $\mathcal{T}(\varepsilon)$ . The different signs are indicated in Fig. 8: the left side of the peaks correspond to  $\Delta\mu < 0$  and the right side to  $\Delta\mu > 0$ . As already mentioned in Ref. 46, these values of  $ZT$  are remarkably high.

Note that the values of  $\mu/(\hbar e/2m^*)$  associated to (a), (b) and (c) do not correspond to maximum values of  $ZT$ , while they are perfectly consistent with the results shown in Fig. 7. Interestingly, the values of the COP calculated from  $ZT$  of Fig. 10 are the same as the maximum ones in Fig. 7 for  $\Delta T/T_h = 0.1$ . This means that the results of Fig. 7 are very close to the conditions for linear response operation.

## V. SUMMARY AND CONCLUSIONS

We have used experimental data of the conductance of the Corbino device studied in Ref. 46 to infer the transmission function. In that reference, it was experimentally verified that the thermoelectric transport takes place in the diffusive regime, without effects introduced by the electron-phonon coupling, implying that it can be accurately described in terms of the transmission function so calculated. Here, we have used that transmission

function to evaluate the cooling power and the COP for cooling.

According to our analysis, in the present sample, the cooling power relative to  $J_{qb}$  for this system is maximum in the vicinity of filling factor  $\nu = 8$ , and the resulting value is:  $(J_c/J_{qb})^{max} = 0.346$  at  $\mu/(\hbar e/2m^*) = 1.85$ ,  $T_h = 170$  mK, for a thermal load  $\Delta T = 5.5$  mK. The chemical potential difference associated to this case results in a bias voltage of  $V = -60$   $\mu$ V. We found that the temperature difference can be increased without a significant decrease in the cooling power.

We have also analyzed the cooling coefficient of performance, as well as the figure of merit characterizing the performance in linear response and we found interesting results, corresponding to large COP ( $\sim 0.3\eta_C$  and larger) for several filling factors, mainly  $\nu = 4, 6, 8$  in the subkelvin regime at a sizable cooling power.

Our estimates, did not include the phononic contribution of the thermal conductance. According to Ref. 44, the latter follows a law  $\kappa_{ph} = \alpha_{ph} T^3$ . We expect the thermoelectric cooling operation to be particularly useful at low temperatures when the electronic cooling power  $J_c \propto T^2$  should overcome the phononic one, induced by the thermal bias.

The specific values obtained from our analysis are associated to the specific sample and device of Ref. 46. However, we have verified that the orders of magnitude of the different quantities are valid for a wide range of samples with similar mobility and devices with similar dimensions, including samples doped with Cr, in which case the phonon conductance is expected to be reduced [50]. Therefore, the present results constitute a proof of principle for the applicability of Landau cooling with Corbino devices in cryogenic conditions. On the other hand, once the transmission function is inferred, the roadmap of our analysis, based on the results presented in Ref. 35, can be also implemented in any other similar particular realization.

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[1] Giuliano Benenti, Giulio Casati, Keiji Saito, and Robert S. Whitney. Fundamental aspects of steady-state

conversion of heat to work at the nanoscale. *Phys. Rep.*, 694:1–124, Jun 2017.

- [2] Francesco Giazotto, Tero T. Heikkilä, Arttu Luukanen, Alexander M. Savin, and Jukka P. Pekola. Opportunities for mesoscopies in thermometry and refrigeration: Physics and applications. *Rev. Mod. Phys.*, 78(1):217–274, Mar 2006.
- [3] A. Ozaeta, P. Virtanen, F. S. Bergeret, and T. T. Heikkilä. Predicted very large thermoelectric effect in ferromagnet-superconductor junctions in the presence of a spin-splitting magnetic field. *Phys. Rev. Lett.*, 112:057001, Feb 2014.
- [4] G. Marchegiani, A. Braggio, and F. Giazotto. Nonlinear thermoelectricity with electron-hole symmetric systems. *Phys. Rev. Lett.*, 124:106801, Mar 2020.
- [5] A.Yu. Kitaev. Fault-tolerant quantum computation by anyons. *Annals of Physics*, 303:2–30, 2003.
- [6] Chetan Nayak, Steven H. Simon, Ady Stern, Michael Freedman, and Sankar Das Sarma. Non-abelian anyons and topological quantum computation. *Rev. Mod. Phys.*, 80:1083–1159, Sep 2008.
- [7] J. R. Prance, C. G. Smith, J. P. Griffiths, S. J. Chorley, D. Anderson, G. A. C. Jones, I. Farrer, and D. A. Ritchie. Electronic Refrigeration of a Two-Dimensional Electron Gas. *Phys. Rev. Lett.*, 102(14):146602, Apr 2009.
- [8] Yanchao Zhang, Guoxing Lin, and Jincan Chen. Three-terminal quantum-dot refrigerators. *Phys. Rev. E*, 91(5):052118, May 2015.
- [9] O. Entin-Wohlman, Y. Imry, and A. Aharony. Enhanced performance of joint cooling and energy production. *Phys. Rev. B*, 91(5):054302, Feb 2015.
- [10] Pablo Roura-Bas, Liliana Arrachea, and Eduardo Fradkin. Helical spin thermoelectrics controlled by a side-coupled magnetic quantum dot in the quantum spin Hall state. *Phys. Rev. B*, 98(19):195429, Nov 2018.
- [11] Davide Venturelli, Rosario Fazio, and Vittorio Giovannetti. Minimal Self-Contained Quantum Refrigeration Machine Based on Four Quantum Dots. *Phys. Rev. Lett.*, 110(25):256801, Jun 2013.
- [12] Jens Schulenburg, Angelo Di Marco, Joren Vanherck, Maarten R. Wegewijs, and Janine Splettstoesser. Thermoelectrics of Interacting Nanosystems—Exploiting Superselection Instead of Time-Reversal Symmetry. *Entropy*, 19(12):668, Dec 2017.
- [13] Rafael Sánchez, Pablo Burset, and Alfredo Levy Yeyati. Cooling by Cooper pair splitting. *Phys. Rev. B*, 98(24):241414(R), Dec 2018.
- [14] Bhaskaran Muralidharan and Milena Grifoni. Performance analysis of an interacting quantum dot thermoelectric setup. *Phys. Rev. B*, 85:155423, Apr 2012.
- [15] D. B. Karki and Mikhail N. Kiselev. Effects of strong electron interactions and resonant scattering on power output of nano-devices. *Phys. Rev. B*, 100:195425, Nov 2019.
- [16] Diego Pérez Daroca, Pablo Roura-Bas, and Armando A. Aligia. Enhancing the nonlinear thermoelectric response of a correlated quantum dot in the kondo regime by asymmetrical coupling to the leads. *Phys. Rev. B*, 97:165433, Apr 2018.
- [17] S. Roddaro, D. Ercolani, M. A. Safeen, S. Suomalainen, F. Rossella, F. Giazotto, L. Sorba, and F. Beltram. Giant thermovoltage in single InAs nanowire field-effect transistors. *Nano Letters*, 13(8):3638–3642, July 2013.
- [18] Prete D., Erdman P. A., Demontis V., Zannier V., Ercolani D., Sorba L., Beltram F., Rossella F., Taddei F., and Roddaro S. Thermoelectric conversion at 30 k in InAs/InP nanowire quantum dots. *Nano Letters*, 19(5):3033–3039, April 2019.
- [19] H. van Houten, L. W. Molenkamp, C. W. J. Beenakker, and C. T. Foxon. Thermo-electric properties of quantum point contacts. *Semicond. Sci. Technol.*, 7(3B):B215–B221, Mar 1992.
- [20] Robert S. Whitney. Nonlinear thermoelectricity in point contacts at pinch off: A catastrophe aids cooling. *Phys. Rev. B*, 88(6):064302, Aug 2013.
- [21] Robert S. Whitney. Finding the quantum thermoelectric with maximal efficiency and minimal entropy production at given power output. *Phys. Rev. B*, 91(11):115425, Mar 2015.
- [22] Kaoru Yamamoto, Amnon Aharony, Ora Entin-Wohlman, and Naomichi Hatano. Thermoelectricity near anderson localization transitions. *Phys. Rev. B*, 96:155201, Oct 2017.
- [23] D. G. Rothe, E. M. Hankiewicz, B. Trauzettel, and M. Guigou. Spin-dependent thermoelectric transport in hgte/cdte quantum wells. *Phys. Rev. B*, 86:165434, Oct 2012.
- [24] Sun-Yong Hwang, Rosa López, Minchul Lee, and David Sánchez. Nonlinear spin-thermoelectric transport in two-dimensional topological insulators. *Phys. Rev. B*, 90(11):115301, Sep 2014.
- [25] Rafael Sánchez, Björn Sothmann, and Andrew N. Jordan. Chiral Thermoelectrics with Quantum Hall Edge States. *Phys. Rev. Lett.*, 114(14):146801, Apr 2015.
- [26] Patrick P. Hofer and Björn Sothmann. Quantum heat engines based on electronic mach-zehnder interferometers. *Phys. Rev. B*, 91:195406, May 2015.
- [27] Luca Vannucci, Flavio Ronetti, Giacomo Dolcetto, Matteo Carrega, and Maura Sasseti. Interference-induced thermoelectric switching and heat rectification in quantum Hall junctions. *Phys. Rev. B*, 92(7):075446, Aug 2015.
- [28] Flavio Ronetti, Luca Vannucci, Giacomo Dolcetto, Matteo Carrega, and Maura Sasseti. Spin-thermoelectric transport induced by interactions and spin-flip processes in two-dimensional topological insulators. *Phys. Rev. B*, 93:165414, Apr 2016.
- [29] Pablo Roura-Bas, Liliana Arrachea, and Eduardo Fradkin. Enhanced thermoelectric response in the fractional quantum Hall effect. *Phys. Rev. B*, 97(8):081104(R), Feb 2018.
- [30] S. Böhlting, G. Engelhardt, G. Platero, and G. Schaller. Thermoelectric performance of topological boundary modes. *Phys. Rev. B*, 98:035132, Jul 2018.
- [31] Sara Kheradsoud, Nastaran Dashti, Maciej Misiorny, Patrick P. Potts, Janine Splettstoesser, and Peter Samuelsson. Power, Efficiency and Fluctuations in a Quantum Point Contact as Steady-State Thermoelectric Heat Engine. *Entropy*, 21(8):777, Aug 2019.
- [32] David Sánchez, Rafael Sánchez, Rosa López, and Björn Sothmann. Nonlinear chiral refrigerators. *Phys. Rev. B*, 99(24):245304, Jun 2019.
- [33] Daniel Gresta, Mariano Real, and Liliana Arrachea. Optimal Thermoelectricity with Quantum Spin Hall Edge States. *Phys. Rev. Lett.*, 123(18):186801, Oct 2019.
- [34] Gianmichele Blasi, Fabio Taddei, Liliana Arrachea, Matteo Carrega, and Alessandro Braggio. Nonlocal thermoelectricity in a superconductor-topological-insulator-superconductor junction in contact with a normal-metal probe: Evidence for helical edge states. *Phys. Rev. Lett.*,

- 124:227701, Jun 2020.
- [35] Fatemeh Hajiloo, Pablo Terrén Alonso, Nastaran Dashti, Liliana Arrachea, and Janine Splettstoesser. Detailed study of nonlinear cooling with two-terminal configurations of topological edge states. *Phys. Rev. B*, 102:155434, Oct 2020.
- [36] L. D. Hicks and M. S. Dresselhaus. Effect of quantum-well structures on the thermoelectric figure of merit. *Phys. Rev. B*, 47(19):12727–12731, May 1993.
- [37] L. D. Hicks and M. S. Dresselhaus. Thermoelectric figure of merit of a one-dimensional conductor. *Phys. Rev. B*, 47(24):16631(R)–16634, Jun 1993.
- [38] G. D. Mahan and J. O. Sofo. The best thermoelectric. *Proc. Natl. Acad. Sci. U.S.A.*, 93(15):7436–7439, Jul 1996.
- [39] Joseph P. Heremans, Mildred S. Dresselhaus, Lon E. Bell, and Donald T. Morelli. When thermoelectrics reached the nanoscale. *Nat. Nanotechnol.*, 8(7):471–473, Jun 2013.
- [40] Yafis Barlas and Kun Yang. Thermopower of quantum hall states in corbino geometry as a measure of quasiparticle entropy. *Phys. Rev. B*, 85:195107, May 2012.
- [41] N. d’Ambrumenil and R. H. Morf. Thermopower in the quantum hall regime. *Phys. Rev. Lett.*, 111:136805, Sep 2013.
- [42] H. van Zalinge, R. W. van der Heijden, and J. H. Wolter. Anisotropic corbino magnetothermopower in a quantum hall system. *Phys. Rev. B*, 67:165311, Apr 2003.
- [43] Shuhei Kobayakawa, Akira Endo, and Yasuhiro Iye. Diffusion thermopower of quantum hall states measured in corbino geometry. *Journal of the Physical Society of Japan*, 82(5):053702, 2013.
- [44] W. E. Chickering, J. P. Eisenstein, L. N. Pfeiffer, and K. W. West. Thermopower of two-dimensional electrons at filling factors  $\nu = 3/2$  and  $5/2$ . *Phys. Rev. B*, 81:245319, Jun 2010.
- [45] W. E. Chickering, J. P. Eisenstein, L. N. Pfeiffer, and K. W. West. Thermoelectric response of fractional quantized hall and reentrant insulating states in the  $n = 1$  landau level. *Phys. Rev. B*, 87:075302, Feb 2013.
- [46] Mariano Real, Daniel Gresta, Christian Reichl, Jürgen Weis, Alejandra Tonina, Paula Giudici, Liliana Arrachea, Werner Wegscheider, and Werner Dietsche. Thermoelectricity in quantum hall corbino structures. *Physical Review Applied*, 14(3):034019, 2020.
- [47] Peter Samuelsson, Sara Kheradsoud, and Björn Sothmann. Optimal Quantum Interference Thermoelectric Heat Engine with Edge States. *Phys. Rev. Lett.*, 118(25):256801, Jun 2017.
- [48] Francesco Giazotto, Fabio Taddei, Michele Governale, Rosario Fazio, and Fabio Beltram. Landau cooling in metal–semiconductor nanostructures. *New Journal of Physics*, 9(12):439–439, dec 2007.
- [49] J. B. Pendry. Quantum limits to the flow of information and entropy. *J. Phys. A: Math. Gen.*, 16(10):2161–2171, Jul 1983.
- [50] N Chaudhuri, RS Wadhwa, Phoola Tiku, and AK Sreedhar. Thermal conductivity of gallium arsenide at low temperatures. *Physical Review B*, 8(10):4668, 1973.